such a deformation of k = 483 N m⁻¹ for 1 and K = 640 N m⁻¹ for 1⁺. This result contradicts what is observed (and by far!) if one assumes that only totally symmetric vibrational modes are stimulated by electron ejection from $3a_1'$. One, as yet unexplored, explanation for the observed, close spacing of the vibrational fine structure of band 1 could be looked for in the vibronic mixing of the ${}^{2}A_{1}$ ground state of 1⁺ with the first excited ${}^{2}E''$ state, which are only $\sim 1.0 \text{ eV}$ apart (adiabatically).

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Registry No. 1, 35634-10-7.

Rotational Spectra and Structures of Small Clusters: Ar₃-HF and Ar₃-DF

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Previous studies of microwave rotational spectra for van der Waals and hydrogen bonded complexes have been limited to binary species like Ar-HF^{1,2} and (HF)₂.³⁻⁵ Herein we demonstrate the feasibility of extending such studies to small molecular clusters. In particular, we report the characterization of Ar₃-HF and Ar₃-DF, in which the H/DF lies along the threefold (c) axis of the Ar₃ group, with the H/D end closest to Ar₃. The symmetric top transitions observed for $J = 1 \rightarrow 2$ through $J = 6 \rightarrow 7$ are fitted by rotational constants B_0 , D_J , and D_{JK} of 1188.212 MHz, 6.85 kHz, and -5.76 kHz for Ar₃-HF and of 1180.379 MHz, 6.57 kHz, and -5.15 kHz for Ar₃-DF. The composition of the clusters is established by hyperfine structure of the transitions, by the absence of K states other than 0, ± 3 , and ± 6 , and by the values found for B_0 .

The spectra were observed with the Flygare Mark II spectrometer.⁶ It combines the principles of pulsed Fourier transform spectroscopy, a Fabry-Perot cavity, and synchronization of the microwave pulse with a pulsed beam of a gas mixture from a supersonic nozzle which cools the gas mixture and helps generate the molecular complexes.⁷ Detection of the cluster was made possible by modifications in the spectrometer which increased S/N about 10-fold by better utilization of its inherent sensitivity.⁸ The main improvements were the use of several microwave pulses after each gas pulse, a 10-fold increase in the repetition rate of the gas pulse, and mounting the gas nozzle to enable ready adjustment of its distance from the axis of the Fabry-Perot mirrors. Also, a microwave coupling system was developed to extend the cavity's operating range to C band (4-8 GHz).8

Argon was used as the carrier gas with 0.5-1.0% of HF and/or DF added. At the nozzle the gas mixture was at ambient temperature and 1 atm. The nozzle diameter was 0.94 mm. Under these conditions Ar–HF and $(HF)_2$ are readily observed,¹⁻⁴ and small clusters of $Ar_n^{9,10}$ and $(HF)_n^3$ are generated at concentrations

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Table I, Observed and Calculated Frequencies of the Rotational Transitions for Ar3-HFª

| $J, K \rightarrow J', K'$ | obsd, MHz | calcd, MHz | difference, kHz |
|-----------------------------|-------------------|------------|-----------------|
| 1,0 → 2,0 | 4752.628 | 4752.630 | -2 |
| $2,0 \rightarrow 3,0$ | 7128.533 | 7128.535 | -2 |
| 3,0 → 4,0 | 9503.944 | 9503.946 | -2 |
| $3,\pm3 \rightarrow 4,\pm3$ | 4.37 ^b | 4.362 | |
| $4,0 \rightarrow 5,0$ | 11878.699 | 11878.700 | -1 |
| $4,\pm3 \rightarrow 5,\pm3$ | 9.21 ^b | 9.219 | |
| $5,0 \rightarrow 6,0$ | 14252.636 | 14252.632 | 4 |
| $5,\pm3 \rightarrow 6,\pm3$ | 3.257 | 3.254 | 3 |
| 6,0 → 7,0 | 16625.579 | 16625.578 | 1 |
| $6,\pm3 \rightarrow 7,\pm3$ | 6.300 | 6.303 | -3 |
| $6,\pm6 \rightarrow 7,\pm6$ | 8.479 | 8.479 | 0 |

^a Transition frequencies for states with $K \neq 0$ are emphasized by omitting the first 3 or 4 digits, which are the same as for the transition just above. ^b Approximate center of partially resolved hyperfine structure; not used in determination of rotational constants.

Table II, Observed and Calculated Frequencies of the Hyperfine Components in the $K = 0, J = 1 \rightarrow 2$ Transitions

| components | | obsd. | calcd. | difference, | | | |
|---|---------------|-----------------------------|--------|-------------|-----|--|--|
| J,F_1,F | \rightarrow | J', F_1', F' | MHz | MHz | kHz | | |
| Transition at 4752.628 MHz for HF Species | | | | | | | |
| $1^{3}/_{2},2$ | | $2^{3}/_{2},^{2}$ | 2.597 | 2.597 | 0 | | |
| $1, \frac{1}{2}, 0$ | | $2,\frac{3}{2},\frac{1}{2}$ | 2.603 | 2.602 | 1 | | |
| $1, \frac{1}{2}, 1$ | | $2,\frac{3}{2},2$ | 2.628 | 2.628 | 0 | | |
| $1, \frac{3}{2}, 1$ | | $2,\frac{5}{2},2$ | | 2.628 | 0 | | |
| $1,^{3}/_{2},2$ | | $2,\frac{5}{2},3$ | | 2.630 | -2 | | |
| 1, 1/2, 1 | | $2^{3}/_{2}, 1$ | 2.679 | 2.678 | 1 | | |
| Transition at 4721.304 MHz for DF Species | | | | | | | |
| $1,1,^{3}/_{2}$ | | $2,1,^{3}/_{2}$ | 1.215ª | 1.216 | -1 | | |
| $1, 1, \frac{1}{2}$ | | $2,2,\frac{3}{2}$ | 1.300 | 1.298 | 2 | | |
| $1,2,\frac{3}{2}$ | | 2,3,5/2 | | 1.299 | 1 | | |
| $1,2,\frac{5}{2}$ | | $2,3,^{7}/_{2}$ | | 1.301 | -1 | | |
| $1, 1, \frac{3}{2}$ | | 2,2,5/2 | 1.307 | 1.306 | 1 | | |
| 1,0,1/2 | | 2,1,3/2 | 1.341 | 1.343 | -2 | | |

^a The first 3 digits are omitted.

decreasing with n. Also, mixed species $Ar_n - (HF)_m$ are formed, with $Ar-(HF)_2$ and Ar_2-HF expected to predominate.¹⁰ Indeed, we have observed a large number of transitions attributed tentatively to Ar₂-HF (T shaped), Ar-(HF)₂, and a noncyclic (HF)₃. The most fully characterized species, however, is a symmetric top, which has the more readily identifiable transitions listed in Table I for HF. They were fitted by the usual Hamiltonian including centrifugal distortion.¹¹ An equivalent set of transitions was observed at somewhat lower frequencies with DF.

Although the $J = 0 \rightarrow 1$ transitions are below the present frequency range, the weaker, less extensive hyperfine structure was observed for the K = 0, $J = 1 \rightarrow 2$ transitions² (Table II). It establishes the presence of an H/DF molecule at the figure axis of the symmetric top, without perturbations from other $I \neq 0$ nuclei. The hyperfine structure was fitted as done earlier² for the Ar-HF and Ar-DF dimers. The adjustable parameters are the line center and the "average" torsional angle α between the H/DF axis and the c-inertial axis, which reduces the hyperfine interactions by the factor $\langle 3 \cos^2 \alpha - 1 \rangle / 2$. The values found for α in the HF and DF species are 41.0 (2)° and 35.6 (4)°, compared with 41.1° and 32.7° reported for the dimers.² The similarity indicates the equilibrium configuration has the H/DF axis along the c axis.

Transitions were observed for only two nonzero K states, assigned in Table I as $K = \pm 3$ and $\pm 6.^{12}$ An alternate assignment is as $K = \pm 1$ and ± 2 , with the higher states too weak for detection. This is precluded, however, by the observed partially resolved lines listed as $K = \pm 3$ for $J = 3 \rightarrow 4$ and $J = 4 \rightarrow 5$. They extend over \sim 50 and 20 kHz, respectively, in accord with calculations

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for $K = \pm 3$, while much narrower widths (≤ 5 kHz) are predicted for $K = \pm 1$. For K states other than 0, ± 3 , and ± 6 to be forbidden, the symmetric top must have a 3-fold axis associated with three zero-spin nuclei.¹¹ These criteria are met by the molecular cluster ${}^{40}Ar_3$ -H/DF, which is an oblate symmetric top.

Confirmation of this composition is provided by analysis of B_0 . Neglecting any tilt of the Ar, plane and vibrations other than torsional oscillations of the H/DF, we have¹³

$$I_{bb}(cluster) = (1/2)(1 + \cos^2 \alpha)I_{bb}(H/DF) + I_{bb}(Ar_3) + \mu_c R^2 (1)$$

where R is the distance between the H/DF and Ar₃ centers of mass (c.m.) and $I_{bb}(Ar_3)$ is a function of r, the Ar-Ar distance. For argon dimer, a B_0 of 0.05778 cm⁻¹ has been determined from a high-resolution Raman study.¹⁴ It gives an r_0 of 3.821 Å, which in eq 1 leads to values of 2.785 and 2.479 Å for R in Ar₃-HF and -DF. The smaller R for the DF species show that the H/D end is pointed at the Ar₃. The corresponding H/DF c.m. to Ar distances of 3.553 and 3.525 Å are comparable with the 3.510 and 3.461 Å found in Ar-H/DF.² Also, the H/DF c.m. to Ar line is 38.5° from the c axis. A full account will be presented later, including the results of a search for the isotopic species ${}^{36}\text{Ar}^{40}\text{Ar}_2$ -H/DF to determine accurate bond distances.

Registry No, Ar, 7440-37-1; HF, 7664-39-3; DF, 14333-26-7.

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Synthesis, Trapping, and Spectral Characterization of 1H-Cyclopropa[1]phenanthrene¹

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The cycloproparenes 1, 3, 2, 4 and 3, 5 and their derivatives are surprisingly stable, readily available materials.⁶ By comparison cyclopropa[a]naphthalene (4) explodes upon melting^{7a} and its 1,1-difluoro derivative^{7b} is stable in solution only at temperatures below -30 °C. This decrease in stability perhaps reflects the formally enhanced olefinic character of the bridge bond. 1H-Cyclopropa[l] phenanthrene (5), the most important unknown structural variant within this family, should exemplify further this

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Scheme I



R1=CI; R2=Se(O)Ph

R¹= H; R²= SeMe 8

R¹ = H; R² = Se Me₂ BF₄

 $R^1 = H; R^2 = SeMeCH_2$





Scheme II



phenomenon. We now report syntheses of 5 and show it to be an unusually reactive molecule stable only at low temperatures.



Recent results have shown⁸ that syn-selenoxide elimination from 6 (Scheme I) occurs to give products derived from 7. However, 7 could be neither isolated nor trapped because of facile cleavage of the three-membered ring. Despite its limitations, the study⁸ has provided the first definitive evidence for the existence of the 1H-cyclopropa[l] phenanthrene ring system and implies that an elimination across the 1a,9b-positions of an appropriate 1a,9bdihydro derivative is the method of choice^{6,8,9} for the synthesis of 5.

With this in view, selenide $8^{10,11}$ has been transformed into the selenonium salt 9 which eliminates dimethyl selenide upon

^{(1) (}a) Studies in the Cycloproparene Series. For the previous part, see ref 8; (b) Presented in part at the 7th I.U.P.A.C. Conference on Physical Organic Chemistry, Auckland, 1984; Abstract A28.

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⁽¹¹⁾ All new compounds reported provided satisfactory analytical and spectral data.